

# Entropic Repulsion for Random Interlacements

Xinyi Li, BICMR, Peking University

Joint work with Zijie Zhuang (Penn)

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- In probability theory, this phenomenon is interpreted as follows: in many cases, a typical realization of some statistical physics model is not the configuration with the lowest energy, but rather from some collection of configurations with high cardinality (entropy).
- The fact that BM typically travels a distance of  $O(\sqrt{t})$  in time  $t$  can be regarded as a result of an entropic force.

# Entropic repulsion for Gaussian free fields

- Consider a random walk bridge  $W : [0, \dots, 2N] \rightarrow \mathbb{Z}$  such that  $W[0] = W[2N] = 0$ . If we condition on the (rare) “hard-wall” event  $\{W[k] \geq 0, k = 0, \dots, 2N\}$ , then the mean value of  $W[N]$  would be “pushed” to  $O(\sqrt{N})$  from 0.

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- In fact, the RW bridge itself can be regarded as a 1D Gaussian free field with Dirichlet boundary condition.

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  - Let  $\Omega_N$  stand for the event  $\{X(z) \geq 0, \forall z \in K_N\}$ . Then

$$\lim_{N \rightarrow \infty} \frac{1}{N^{d-2} \log N} \log P[\Omega_N] = c_d \text{cap}_B(K)$$

where  $\text{cap}_B(K)$  stands for the Newtonian capacity of  $K$ .

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- (Deuschel and Giacomin, *CMP* '99) Moreover, as  $N \rightarrow \infty$ ,

$$P\left(X(\cdot) - \sqrt{c'_d \log N} \Big| \Omega_N\right) \Rightarrow P(X(\cdot)).$$

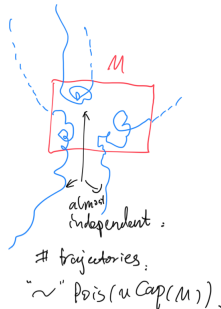
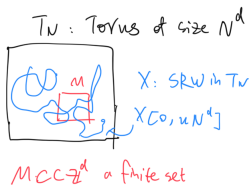
- In other words, the field is “pushed” to infinity by the hard-wall conditioning.
- Similar results also hold for the 2D case, ( $P[\Omega_N] \approx N^{-c \log N}$ , field pushed to  $c_2 \log N$  conditioned on  $\Omega_N$ ), see Bolthausen, Deuschel and Giacomin (*AoP*, '01).

## Random interlacements

- The model of random interlacements was introduced by Sznitman (*Ann. Math.* '10) to describe the local picture left by the trace of a random walk on a large discrete torus ( $d \geq 3$ ) when it runs up to a time proportional to the volume of the torus.

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- It can be regarded as a Poissonian cloud of doubly-infinite trajectories, governed by an intensity parameter  $u > 0$ .



## Random interlacements

- Write  $\mathcal{I}^u \subset \mathbb{Z}^d$  for the union of all trajectories that appears in the cloud. We can then characterize  $\mathcal{I}^u$  as a random subset of  $\mathbb{Z}^d$  through the following relation

$$P[\mathcal{I}^u \cap V = \emptyset] = \exp(-u \text{cap}(V))$$

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- One can also regard random interlacements as a model of percolation with strong correlation.
- In fact,

$$\text{Cov}(\mathbf{1}_{0 \in \mathcal{I}^u}, \mathbf{1}_{x \in \mathcal{I}^u}) \sim c(u)|x|^{2-d}.$$

- The interlacement set  $\mathcal{I}^u$  is always well-connected. It is the *vacant set*  $\mathcal{V}^u := \mathbb{Z}^d \setminus \mathcal{I}^u$  on which the phase transition of percolation is non-trivial.

## Percolation on the vacant set

Theorem (Sznitman '10, Sidoravicius-Sznitman '08, Teixeira '08)

There exists  $u_* \in (0, \infty)$  such that

- for all  $u < u_*$ ,  $\mathcal{V}^u$  has a unique infinite cluster,  $\mathbb{P}_u$ -a.s.;
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- The parallel sharpness of phase transition of the level sets of Gaussian free field has been verified recently by Duminil-Copin, Goswami, Rodriguez, Severo ('20+).

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Theorem (L. and Zhuang, '22+, )

(Large deviations) For a compact regular  $K \subset \mathbb{R}^d$ , let  $K_N$  be its discrete blow-up. Consider two independent interlacements  $\mathcal{I}_1$  and  $\mathcal{I}_2$  with intensities  $u_1, u_2 > 0$  resp.,

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(Large deviations) For compact regular  $K \subset \mathbb{R}^d$  and let  $K_N$  be its discrete blow-up. Consider two independent random interacements  $\mathcal{I}_1$  and  $\mathcal{I}_2$  with intensities  $u_1, u_2 > 0$  resp.

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- The lower bound of the first claim is trivial.
- However, the upper bound requires an involved coarse graining argument.



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Theorem (L. and Zhuang, '22+)

Given  $u > 0$  and consider random interlacements  $\mathcal{I}^u$ , then

$$\lim_{N \rightarrow \infty} \frac{1}{N^{d-2}} \log \mathbb{P}[\text{cap}(K_N \cap \mathcal{I}^u) < \lambda N^{d-2}] = -\frac{u}{d} f_K(d\lambda),$$

## Disconnection by Random Interlacements

- Given  $M > 0$  and a regular  $K \subset [-M, M]^d$  and  $N \in \mathbb{N}$ , look at its discrete blow-up  $K_N$ . And consider the event

$$A_N \triangleq \{K_N \not\leftrightarrow^{\mathcal{V}^u} \partial S_N\} :$$

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- When  $u$  is small such that  $\mathcal{V}^u$  is (strongly) percolative,  $A_n$  is a very unlikely event. A naive bound gives  $P[A_n] \geq c^{-N^{d-1}}$ .

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- Question: What is the precise asymptotics of  $P[A_N]$ ? What is the optimal strategy in realizing  $A_N$ ?
- One can also consider the same type of disconnection by a single trajectory of random walk: let  $X_n$  be a simple random walk on  $\mathbb{Z}^d$ ,  $d \geq 3$ , started from the origin and write  $\mathcal{V} = \mathbb{Z}^d \setminus X[0, \infty)$ . Then define  $A_n^0$  in the same fashion as  $A_N$ , with  $\mathcal{V}^u$  replaced by  $\mathcal{V}$ .



# Disconnection by Random Interlacements and SRW

Theorem (L.-Sznitman *EJP* '14 lower bound, Sznitman *PTRF* '17 upper bound for  $K$  being a box, Nitzschner-Sznitman *JEMS* '20 upper bound for general  $K$ )

For  $0 < u < u_{**}$  and  $0 < u < \bar{u}$  resp.,

$$\liminf_{N \rightarrow \infty} \frac{1}{N^{d-2}} \log P[A_N] \geq -\frac{(\sqrt{u_{**}} - \sqrt{u})^2}{d} \text{cap}_B(K);$$

$$\limsup_{N \rightarrow \infty} \frac{1}{N^{d-2}} \log P[A_N] \leq -\frac{(\sqrt{\bar{u}} - \sqrt{u})^2}{d} \text{cap}_B(K).$$

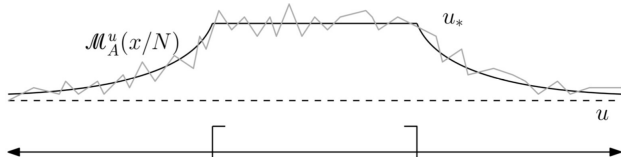
In the case of random walk, similar bounds for  $A_n^0$  hold with  $u = 0$  on the RHS' (the upper bound trivially follows from the RI case, but lower bound (L. AoP '17) requires substantial work and a somewhat different strategy).

## Asymptotics on the Disconnection Probability

- The lower bound is obtained through the introduction of the *tilted interlacements*, i.e., RI defined on  $\mathbb{Z}^d$  with inhomogeneous edge-weights which is modulated by a certain harmonic function.

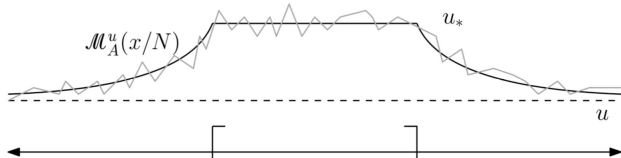
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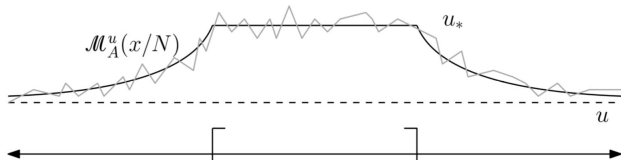
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- The upper bound is obtained through an extremely involved coarse graining argument and capacity estimates.

# Asymptotics on the Disconnection Probability

- The lower bound is obtained through the introduction of the *tilted interlacements*, i.e., RI defined on  $\mathbb{Z}^d$  with inhomogeneous edge-weights which is modulated by a certain harmonic function.



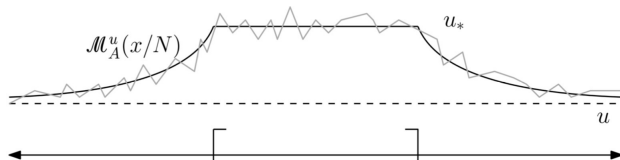
- The upper bound is obtained through an extremely involved coarse graining argument and capacity estimates.
- **If the conjecture that  $\bar{u} = u_* = u_{**}$  is true**, then the upper and lower bounds coincide and we can confirm that the “tilted interlacements” indeed provides an optimal strategy for the disconnection.

# Entropic Repulsion for Random Interlacements Conditioned on disconnection

- Chiarini and Nitzschner (*AoP* '20) showed that **if the conjecture**

$$\bar{u} = u_* = u_{**}$$

**is indeed true**, then conditioned on the disconnection event  $A_n$ , the occupation time profile of the interlacements indeed roughly coincides with that of the tilted interlacements. In other words, this suggests that conditioning on disconnection “pushes” up the intensity of interlacements from  $u$  to  $u_*$  in the bulk of  $K_N$ .



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$$\lim_{t \rightarrow \infty} \frac{1}{t^{(d-2)/d}} \log P(|W^a(t)| \leq bt) = I_d(a, b),$$

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- For the problem we investigated, we conjecture that tilted interlacements also enters into play in the conditioning.
- Conditional on  $(\mathcal{I}_1 \cap \mathcal{I}_2) \cap K_N = \emptyset$ , when  $N$  tends to  $\infty$ , we expect that the law of  $(\mathcal{I}_1, \mathcal{I}_2)$  should resemble more and more  $(\mathcal{I}^{u_1}, \tilde{\mathcal{I}})$  where  $\tilde{\mathcal{I}}$  represents a kind of tilted interlacements with intensity 0 in  $K_N$  and  $u_2$  at infinity, independent of  $\mathcal{I}^{u_1}$ .

Thanks for your attention!